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2024-2025 Course of Power System Analysis Fall Semester

#### **Grounding regime**

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#### Context

A fundamental element for the choice of grounding regime (i.e., the connection of the star center of the three phase system) in medium- and high-voltage networks is the behaviour of the network itself in the case of a short-circuit to ground.

This abnormal operating condition leads to **significant** overvoltages at nominal frequency (50 Hz), which must be managed in order to operate the system correctly. Therefore, we must study the behavior of an electrical network after a single phase to ground short-circuit with a short-circuit impedance equal to  $\bar{Z}$ .

## Outline

Single phase to ground shortcircuit

Grounding regime in mediumand high-voltage distribution networks

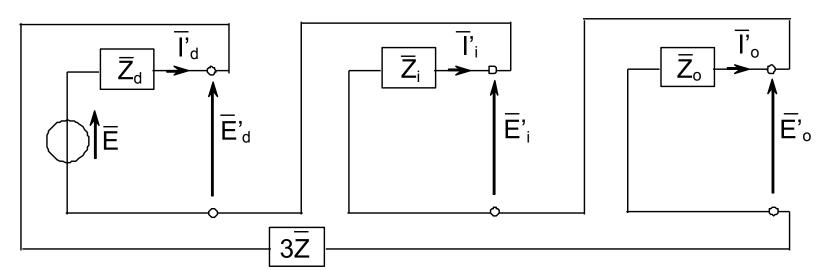
Networks with earthed neutral

Networks with isolated neutral

To recall, the following are the **three equations in the sequence domain** 

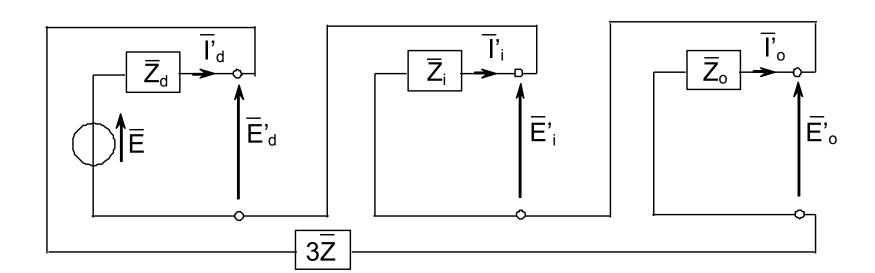
$$\overline{I}_{0}' = \overline{I}_{d}' 
\overline{I}_{d}' = \overline{I}_{i}' 
\overline{E}_{0}' + \overline{E}_{d}' + \overline{E}_{i}' = 3\overline{Z}\overline{I}_{0}'$$

These three equations can be seen as a coupling of the sequence equivalent circuits.



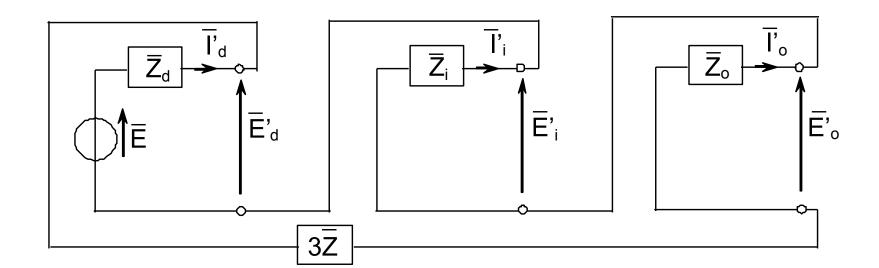
The solution of the circuit enables us to determine the sequence currents:

$$\overline{I}_0' = \overline{I}_d' = \overline{I}_i' = \frac{\overline{E}}{\overline{Z}_d + \overline{Z}_i + \overline{Z}_0 + 3\overline{Z}}$$



Therefore, the **short-circuit current in the phase domain** is:

$$\overline{I}_{a}' = \frac{3\overline{E}}{\overline{Z}_{d} + \overline{Z}_{i} + \overline{Z}_{0} + 3\overline{Z}} = \frac{3\overline{E}}{\overline{Z}_{d} + \overline{Z}_{i} + \overline{Z}_{0}}\Big|_{\overline{Z}=0}$$



#### The **sequence voltages** are:

$$\begin{split} \overline{E}_{d}^{'} &= \overline{E} - \overline{Z}_{d} \overline{I}_{d}^{'} \\ \overline{E}_{i}^{'} &= -\overline{Z}_{i} \overline{I}_{i}^{'} \\ \overline{E}_{0}^{'} &= -\overline{Z}_{0} \overline{I}_{0}^{'} \end{split}$$

Plugging in the current equations, we obtain:

$$\begin{split} \overline{E}'_{d} &= -\left(\overline{E}'_{0} + \overline{E}'_{i}\right) \\ \overline{E}'_{i} &= -\overline{Z}_{i} \frac{\overline{E}}{\overline{Z}_{d} + \overline{Z}_{i} + \overline{Z}_{0}} \\ \overline{E}'_{0} &= -\overline{Z}_{0} \frac{\overline{E}}{\overline{Z}_{d} + \overline{Z}_{i} + \overline{Z}_{0}} \end{split}$$

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Thanks to the transformation matrix for symmetrical components:

$$\begin{bmatrix} \overline{E}'_a \\ \overline{E}'_b \\ \overline{E}'_c \end{bmatrix} = \begin{bmatrix} 1 & 1 & 1 \\ 1 & \alpha^2 & \alpha \\ 1 & \alpha & \alpha^2 \end{bmatrix} \begin{bmatrix} \overline{E}'_0 \\ \overline{E}'_d \\ \overline{E}'_i \end{bmatrix}$$

It is possible to express the voltages in the phase domain.

$$\overline{E}_{a}' = \overline{E}_{0}' + \overline{E}_{d}' + \overline{E}_{i}' = 0$$

$$\overline{\overline{E}}_{b}' = \overline{\overline{E}}_{0}' + \alpha^{2} \overline{\overline{E}}_{d}' + \alpha \overline{\overline{E}}_{i}' = -\overline{Z}_{0} \frac{\overline{\overline{E}}}{\overline{Z}_{d} + \overline{Z}_{i} + \overline{Z}_{0}} - \alpha^{2} \left(\overline{\overline{E}}_{0}' + \overline{\overline{E}}_{i}'\right) - \alpha \overline{Z}_{i} \frac{\overline{\overline{E}}}{\overline{Z}_{d} + \overline{Z}_{i} + \overline{Z}_{0}}$$

$$=\frac{\overline{E}}{\overline{Z}_d+\overline{Z}_i+\overline{Z}_0}\left[-\overline{Z}_0-\alpha^2\left(\overline{E}_0'+\overline{E}_i'\right)\frac{\overline{\overline{Z}_d+\overline{Z}_i+\overline{Z}_0}}{\overline{E}}-\alpha\overline{Z}_i\right]=$$

$$=\frac{\overline{E}}{\overline{Z}_d+\overline{Z}_i+\overline{Z}_0}\begin{bmatrix} -\overline{Z}_0-\alpha^2\frac{\overline{(\overline{Z}_0+\overline{Z}_i)}}{\overline{I}_0}-\alpha\overline{Z}_i \end{bmatrix}=$$

$$\frac{\overline{E}}{\overline{Z}_{d} + \overline{Z}_{i} + \overline{Z}_{0}} \left[ -\overline{Z}_{0} + \alpha^{2} \overline{Z}_{0} + \alpha^{2} \overline{Z}_{i} - \alpha \overline{Z}_{i} \right] = \Big|_{\overline{E} = E + j0} E \frac{\left(\alpha^{2} - 1\right) \overline{Z}_{0} + \left(\alpha^{2} - \alpha\right) \overline{Z}_{i}}{\overline{Z}_{d} + \overline{Z}_{i} + \overline{Z}_{0}}$$

$$E \frac{\left(\alpha^2 - 1\right)\overline{Z}_0 + \left(\alpha^2 - \alpha\right)\overline{Z}_i}{\overline{Z}_d + \overline{Z}_i + \overline{Z}_0}$$

$$\begin{split} & \overline{E}_{c}' = \overline{E}_{0}' + \alpha \overline{E}_{d}' + \alpha^{2} \overline{E}_{i}' = -\overline{Z}_{0} \frac{\overline{E}}{\overline{Z}_{d} + \overline{Z}_{i} + \overline{Z}_{0}} - \alpha \left(\overline{E}_{0}' + \overline{E}_{i}'\right) - \alpha^{2} \overline{Z}_{i} \frac{\overline{E}}{\overline{Z}_{d} + \overline{Z}_{i} + \overline{Z}_{0}} = \\ & = \frac{\overline{E}}{\overline{Z}_{d} + \overline{Z}_{i} + \overline{Z}_{0}} \left[ -\overline{Z}_{0} - \alpha \left(\overline{E}_{0}' + \overline{E}_{i}'\right) \frac{\overline{Z}_{d} + \overline{Z}_{i} + \overline{Z}_{0}}{\overline{E}} - \alpha^{2} \overline{Z}_{i} \right] = \end{split}$$

$$=\frac{\overline{E}}{\overline{Z}_d+\overline{Z}_i+\overline{Z}_0}\left[-\overline{Z}_0-\alpha\frac{\overline{(\overline{Z}_0+\overline{Z}_i)}}{\overline{I}_0}-\alpha^2\overline{Z}_i\right]=$$

$$\frac{\overline{E}}{\overline{Z}_d + \overline{Z}_i + \overline{Z}_0} \left[ -\overline{Z}_0 + \alpha \overline{Z}_0 + \alpha \overline{Z}_i - \alpha^2 \overline{Z}_i \right] = \Big|_{\overline{E} = E + j0} E \frac{\left(\alpha - 1\right) \overline{Z}_0 + \left(\alpha - \alpha^2\right) \overline{Z}_i}{\overline{Z}_d + \overline{Z}_i + \overline{Z}_0}$$

$$E^{\frac{(\alpha-1)\overline{Z}_0 + (\alpha-\alpha^2)\overline{Z}_i}{\overline{Z}_d + \overline{Z}_i + \overline{Z}_0}}$$

We can introduce the ratios of the homopolar and inverse sequence impedances relative to the direct sequence impedance:

$$ar{m}_0 = rac{ar{Z}_0}{ar{Z}_d}; \quad ar{m}_i = rac{ar{Z}_i}{ar{Z}_d}$$

If we substitute these ratios in the previous equations where the phase voltages (b and c) are expressed relative to the voltage E, we obtain:

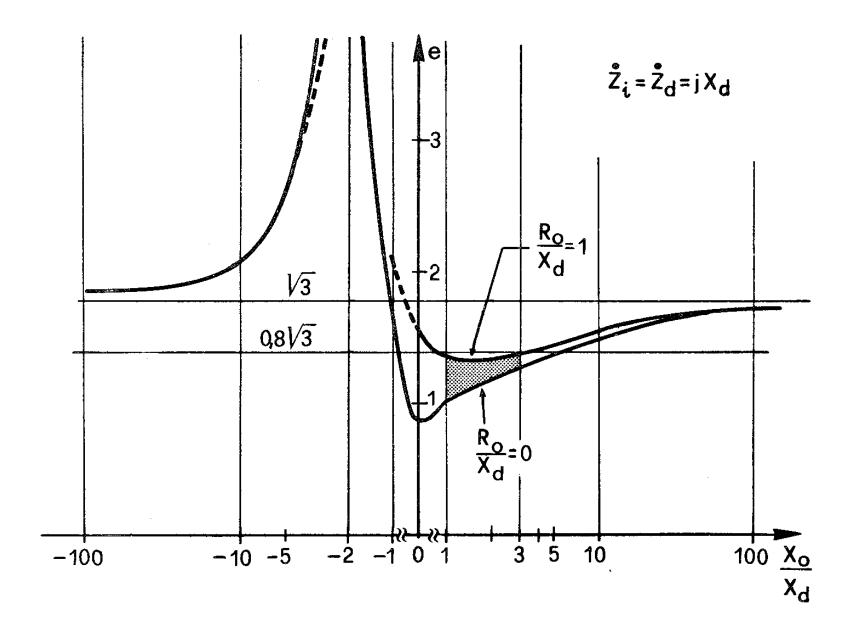
$$\overline{e}_{b}' = \frac{\overline{E}_{b}'}{E} = \frac{\left(\alpha^{2} - 1\right)\overline{m}_{0} + \left(\alpha^{2} - \alpha\right)\overline{m}_{i}}{1 + \overline{m}_{i} + \overline{m}_{0}}$$

$$\overline{e}_{c}' = \frac{\overline{E}_{c}'}{E} = \frac{\left(\alpha - 1\right)\overline{m}_{0} + \left(\alpha - \alpha^{2}\right)\overline{m}_{i}}{1 + \overline{m}_{0} + \overline{m}_{0}}$$

If we consider that the inverse sequence impedance is almost the same as the direct sequence impedance, we can say that  $m_i=1$  and simplify the equations.

Further, if we **neglegt the resistive components** of these impedances in the direct sequence impedance of the network, we have  $\overline{e}_b$  and  $\overline{e}_c$  depending only on the coefficient  $m_0$ , which is equal to

$$\overline{m}_0 = \frac{\overline{Z}_0}{jX_d}$$



Two possible cases:

1. Network with isolated neutral: in these networks, the only possibility for short-circuit current to circulate is provided by the shunt capacitances of the conductors to ground; therefore

$$\bar{Z}_0 = -jX_0$$

where  $X_0$  is a purely capacitive reactance linked to the capacity of the lines in service

$$X_0 = \frac{1}{\omega C_{\text{serv linea}}}$$

Therefore, the larger the network, the smaller the value of  $X_0$ .

Two possible cases:

2. Network with grounded neutral: in these networks, the star center of the system is grounded in the primary substation by an ohmic-inductive impedance, and therefore:

$$\overline{Z}_0 = R_0 + jX_0$$

It is important to observe that the value of  $Z_0$  strongly depends on the presence of inductances connected between the ground and the star center (Petersen winding), and on the behavior of the transformers and lines in the homopolar sequence.

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#### Networks with earthed neutral

A network is defined with **neutral effectively grounded** if, **in** the case of a single phase to ground short-circuit, the ratio k, at any point in the network, between the largest voltage to earth of healthy phases and the phase-to-phase voltage without the short-circuit is no greater than 0.8:

$$e \le 0.8\sqrt{3} = 1.37$$

This ratio k is called the **grounding coefficient** and is one of the parameters for **insulation coordination**. In fact, **short-circuit to ground overvoltages are generally the most important**, and only sometimes can they be exceeded by overvoltages due to sudden tripping of load.

#### Networks with earthed neutral

The determination of the coefficient k, and therefore the prediction of the most important overvoltages, is used to establish whether the 50 Hz test voltage set point for machines and devices is properly chosen.

The dotted area in the previous figure means that k < 0.8 (neutral effectively grounded) if:

$$R_0 / X_d \le 1 \qquad 1 \le X_0 / X_d \le 3$$

To understand these conditions, we should consider two cases:

- Transmission networks
- Distribution networks

#### Networks with earthed neutral

**Transmission networks**: to obtain a sufficiently low impedance  $Z_0$  in any section of the network, the **star** centers of transformers in substations should be connected to ground.

Distribution networks: as soon as the neutral is connected to ground at a single point (medium-voltage side of the primary substation transformer), a neutral to ground connection impedance must be provided by means of an inductance (since the network impedance of the direct sequence  $X_d$  will surely be inductive and the ratio  $X_0/X_d$  must be positive).

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Medium-voltage networks with isolated neutral exhibit, in general, a ratio  $X_0/X_d=-20$  because the homopolar sequence reactance  $X_0$ , represented by the shunt capacitances of the line conductors, is much larger than the direct sequence reactance of the system. For these networks, the previously shown figure shows how the phase voltage of the healthy phases increases by a factor  $\sqrt{3}$ .

Therefore, in a distribution system with isolated neutral, in a case of short-circuit to ground, the single-phase voltage of the healthy phases becomes equal to the phase-to-phase voltage and the short-circuit current has a small value. In effect, the only link between the electrical system and ground is given by the shunt capacitances of each phase.

The shunt capacitances, which have a small value, are the homopolar impedance of the network, which has a high value. Therefore, we can say that

$$\left| \overline{Z}_{0} \right| >> \left| \overline{Z}_{d} \right|$$

And, due to the following relations:

$$\overline{Z}_{d} = \overline{Z}_{i} \qquad |\overline{Z}_{0}| >> |\overline{Z}_{i}|$$

#### The voltages of the healthy phases are:

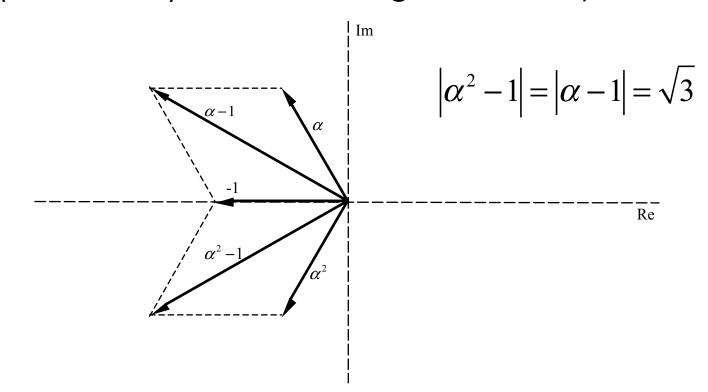
$$\overline{E}'_{b} = E \frac{\left(\alpha^{2} - 1\right)\overline{Z}_{0} + \left(\alpha^{2} - \alpha\right)\overline{Z}_{i}}{\overline{Z}_{d} + \overline{Z}_{i} + \overline{Z}_{0}} \cong E \frac{\left(\alpha^{2} - 1\right)\overline{Z}_{0}}{\overline{Z}_{0}} = E\left(\alpha^{2} - 1\right)$$

$$\overline{E}'_{c} = E \frac{\left(\alpha - 1\right)\overline{Z}_{0} + \left(\alpha - \alpha^{2}\right)\overline{Z}_{i}}{\overline{Z}_{d} + \overline{Z}_{i} + \overline{Z}_{0}} \cong E \frac{\left(\alpha - 1\right)\overline{Z}_{0}}{\overline{Z}_{0}} = E\left(\alpha - 1\right)$$

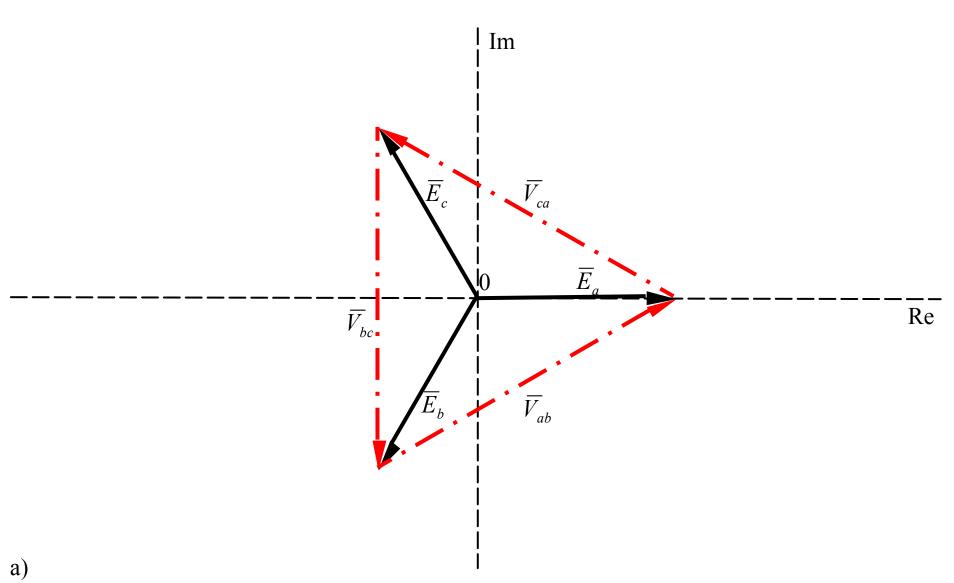
It is interesting to note that the coefficients

$$(\alpha^2-1)$$
  $(\alpha-1)$ 

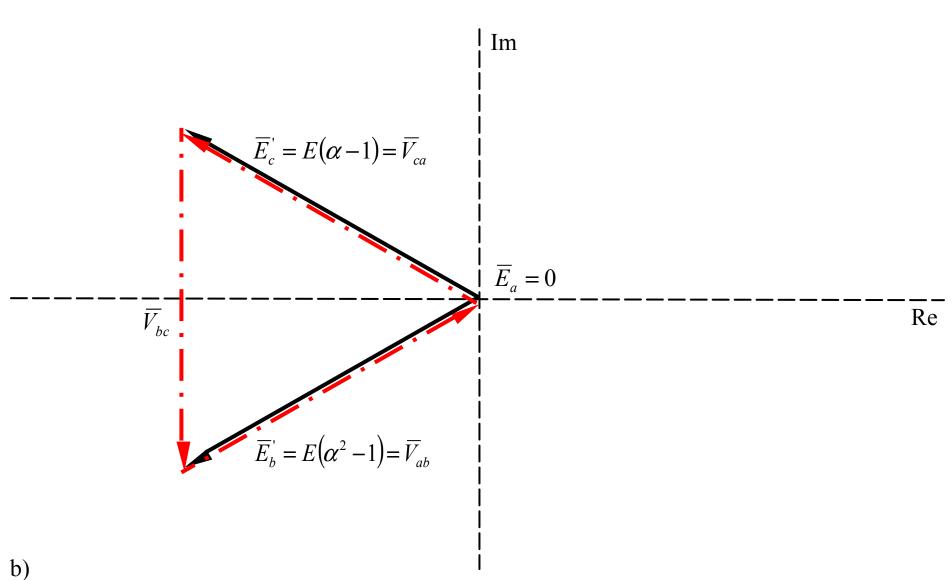
Are two well-defined complex numbers in the Gaussian plane (see below) and their magnitude is equal to:



Before the short circuit:



After the short circuit:



#### **Observation**:

$$\overline{E}_b' = E(\alpha^2 - 1)$$

$$\overline{E}_c' = E(\alpha - 1)$$

Figure a) shows the single phase and phase-to-phase voltages in a network with an isolated neutral before a single phase short-circuit to ground. Figure b) shows the single phase and phase-to-phase voltages after the short-circuit. Looking at Figure b), we can note that:

- i) The phase voltages are as given by the above equations
- ii) The phase-to-phase voltages have not changed in value.

This last remark is coherent with the fact that in a system with isolated neutral, the voltages imposed by the system are phase-to-phase voltages.

It is interesting to note that **in systems with isolated neutral**, the short-circuit current for a single phase to ground short-circuit is generally **low** (on the order of tens or hundreds of amps). Using the equation

$$\overline{I}_{a}' = \frac{3\overline{E}}{\overline{Z}_{d} + \overline{Z}_{i} + \overline{Z}_{0} + 3\overline{Z}} = \frac{3\overline{E}}{\overline{Z}_{d} + \overline{Z}_{i} + \overline{Z}_{0}}\Big|_{\overline{Z}=0}$$

we can see that the short-circuit current is inversely proportional to the sum of the sequence impedances of the network.

Given that

$$\overline{Z}_0 = -jX_0 = \frac{1}{\omega C_{\text{serv linee}}}$$

And that the capacitances of the network are on the order of  $10 \cdot 10^{-12}$  F/km for overhead lines and on the order of  $300 \cdot 10^{-12}$  F/km for cables, we can deduce that  $X_0$  is on the order of hundreds or millions of Ohms (and is a function of the line length).

It is equally important to remark that the short-circuit current and the homopolar voltage can be easily placed in the Gaussian plane.

The homopolar voltage is:

$$\overline{E}_0 = \frac{1}{3} \left( \overline{E}_a' + \overline{E}_b' + \overline{E}_c' \right) =$$

$$= \frac{1}{3} \left( \overline{E}_a' + \overline{E}_b' + \overline{E}_c' - \overline{E}_c$$

#### The **homopolar current** is:

$$I_0' = \overline{I}_d' = \overline{I}_i' = \frac{\overline{E}}{\overline{Z}_d + \overline{Z}_i + \overline{Z}_0 + 3\overline{Z}}$$

Given that

$$\left| \overline{Z}_{0} \right| >> \left| \overline{Z}_{d} \right| \qquad \left| \overline{Z}_{0} \right| >> \left| \overline{Z}_{i} \right| \qquad \overline{Z}_{0} = -jX_{0}$$

And that the voltage E is the voltage that supplies the short-circuited phase a on the real axis, we can write that

$$\overline{E} = E + j0$$

and therefore:

$$\overline{I}_a' = 3\overline{I}_0' = 3\overline{I}_d' = 3\overline{I}_i' = \frac{3\overline{E}}{\overline{Z}_d + \overline{Z}_i + \overline{Z}_0} \cong \frac{3E}{-jX_0} = j\frac{3E}{X_0}$$

